

Splines on the Sphere (A View from the Other Hemisphere)

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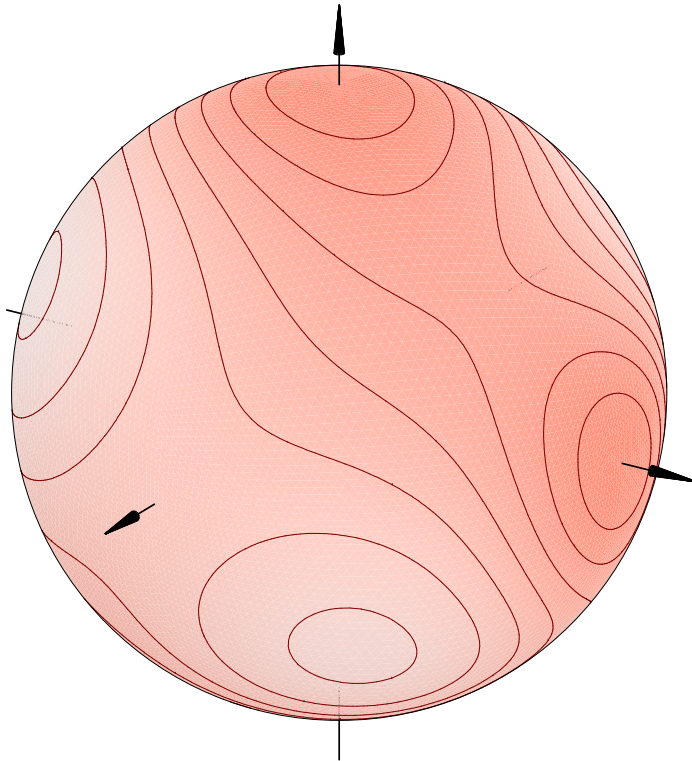
stolfi@ic.unicamp.br

Joint work with

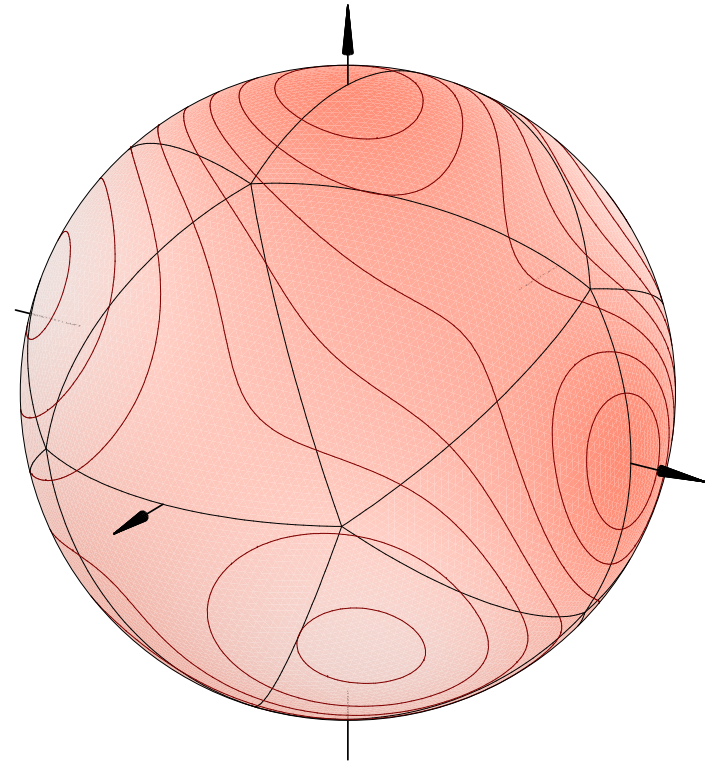
ANAMARIA GOMIDE

Spherical Functions and Spherical Splines

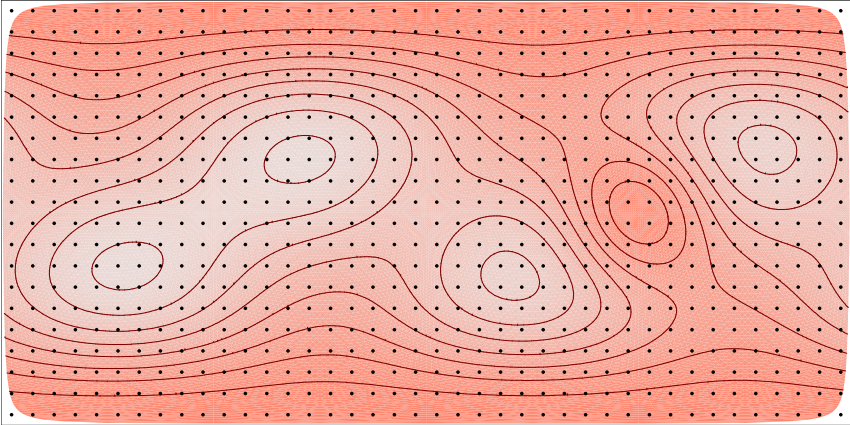
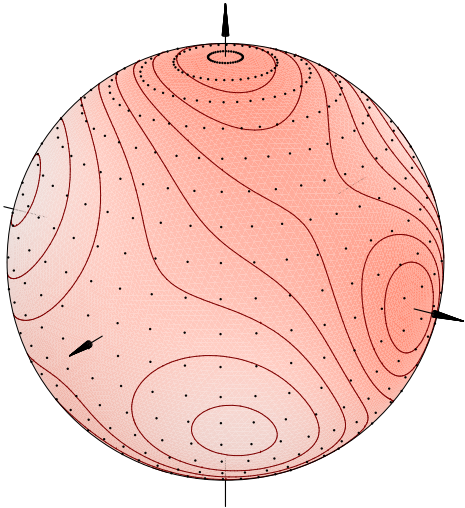
A Spherical function:



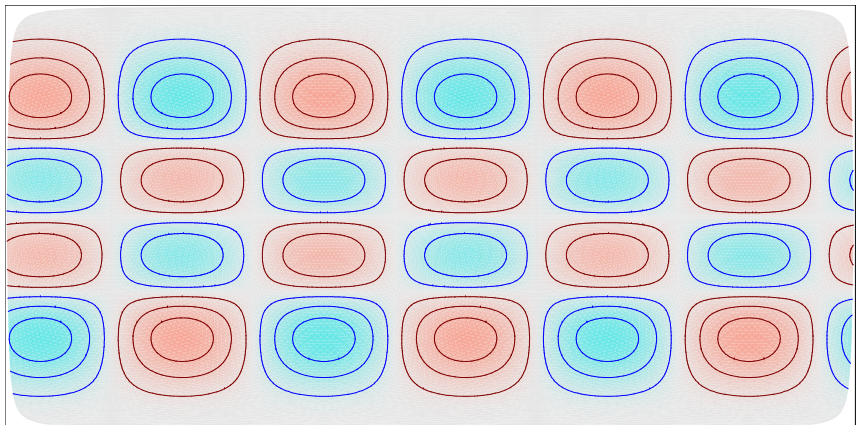
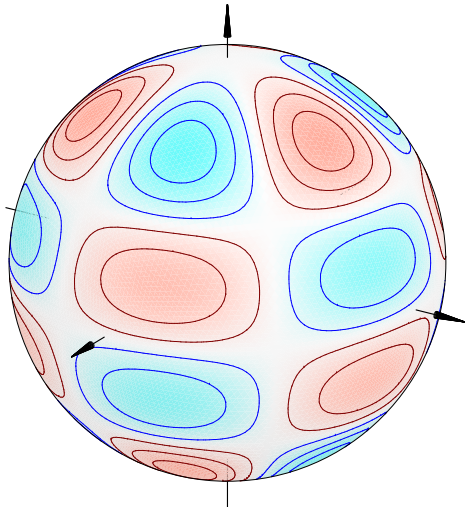
A Spherical spline:



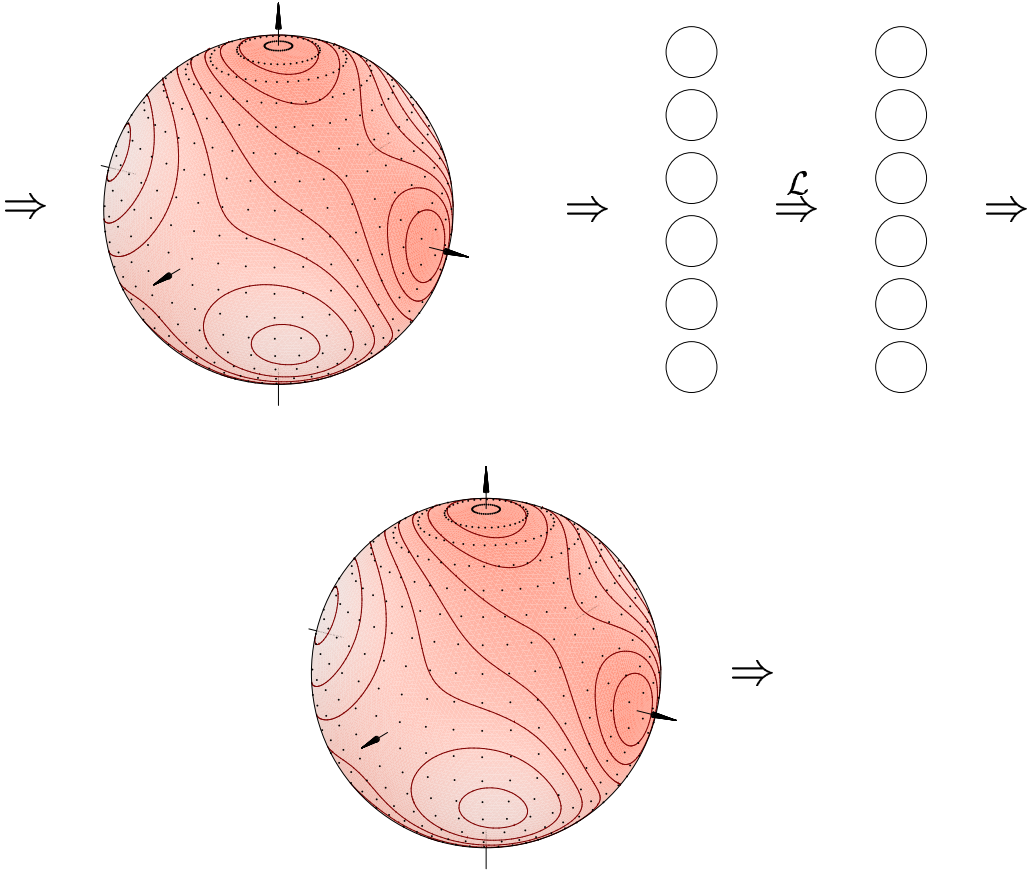
Modeling with Uniform Lon-Lat Grids



Modeling with Spherical Harmonics

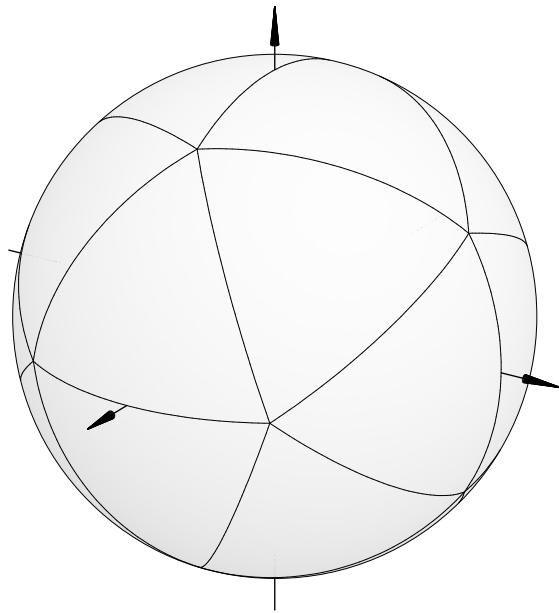


The Spectral Method

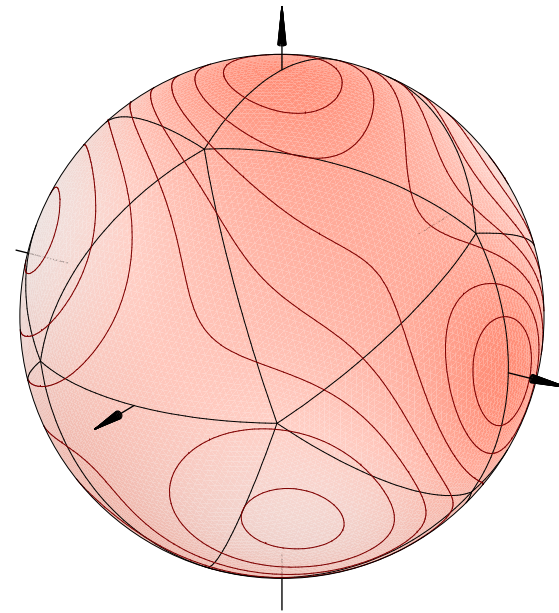


Spherical Splines

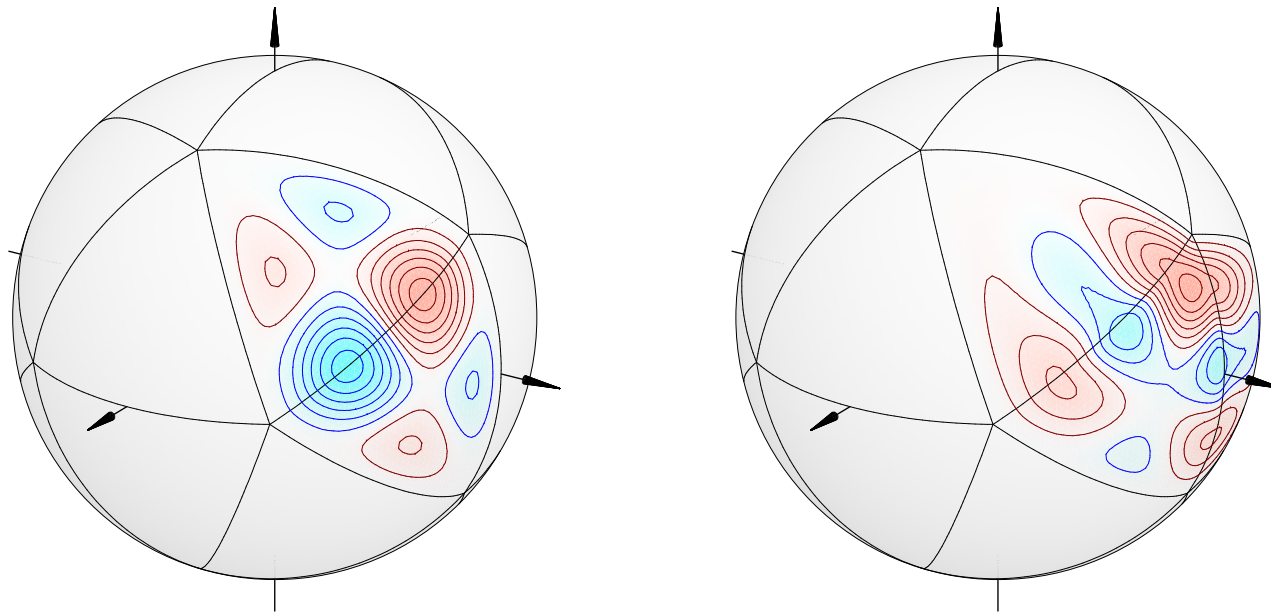
Geodesic triangulation



Spherical spline



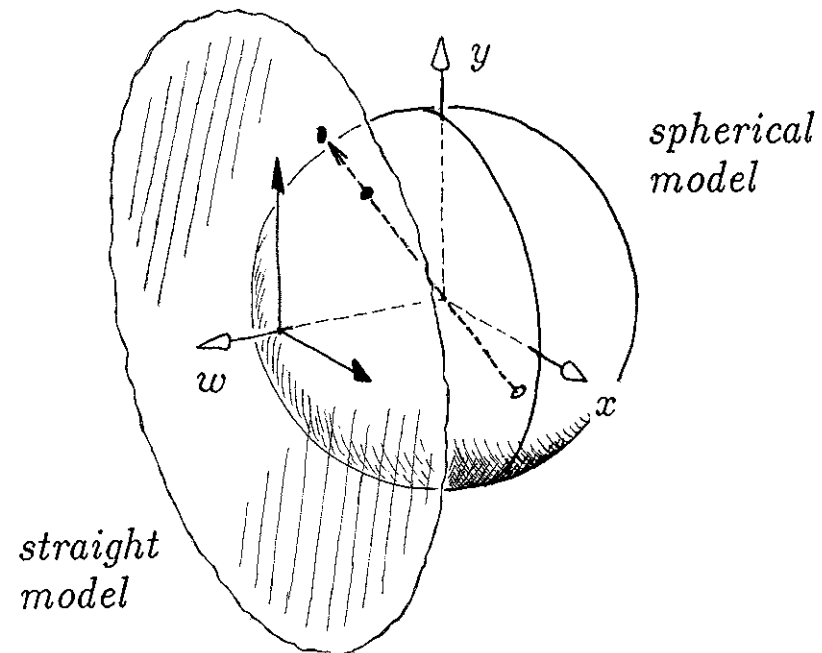
Finite Element Bases for Spherical Splines



Cartesian Approach to Spherical Functions

$$F(x, y, z)$$

$$f = F|_{\mathbf{S}^2}$$



Spherical Polynomials

\mathcal{H}^d = Homogeneous trivariate polynomials of degree d .

\mathcal{P}^d = General trivariate polynomials of degree at most d
= $\mathcal{H}^0 \oplus \mathcal{H}^1 \oplus \dots \oplus \mathcal{H}^d$

On the sphere, $1 = x^2 + y^2 + z^2$.

Therefore $\mathcal{P}^d|_{\mathbf{S}^2} = \mathcal{H}^d|_{\mathbf{S}^2} \oplus \mathcal{H}^{d-1}|_{\mathbf{S}^2}$.

Spherical Splines

$\mathcal{H}_r^d[T]|\mathbf{S}^2$ = Homogeneous C_r spherical splines on T with degree d .

$\mathcal{P}_r^d[T]|\mathbf{S}^2$ = General C_r spherical splines on T with degree at most d
= $\mathcal{H}_r^d[T]|\mathbf{S}^2 \oplus \mathcal{H}_r^{d-1}[T]|\mathbf{S}^2$

[Gomide and Stolfi 1998]

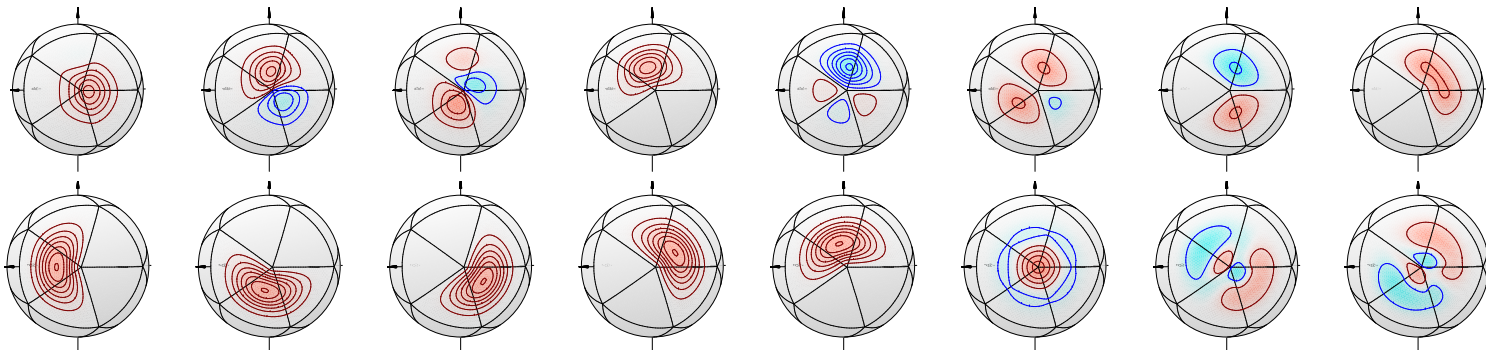
For the circle:

$$\mathcal{P}_r^d[T]|\mathbf{S}^1 \supset \mathcal{H}_r^d[T]|\mathbf{S}^2 \oplus \mathcal{H}_r^{d-1}[T]|\mathbf{S}^2.$$

ANS Finite Element Bases

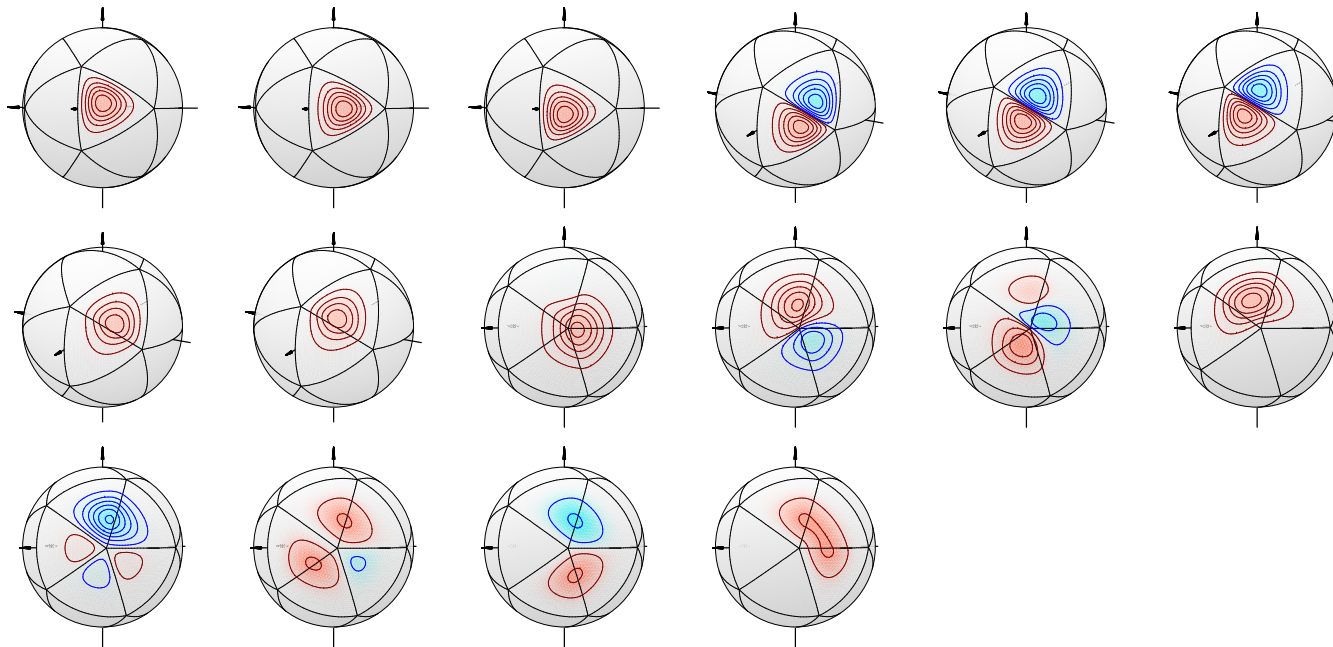
Two bases for each space, differ on vertex elements:

- Original construction (O): $g + 3$ “oranges”.
- New construction (N): g “boats” and 3 “oranges”.



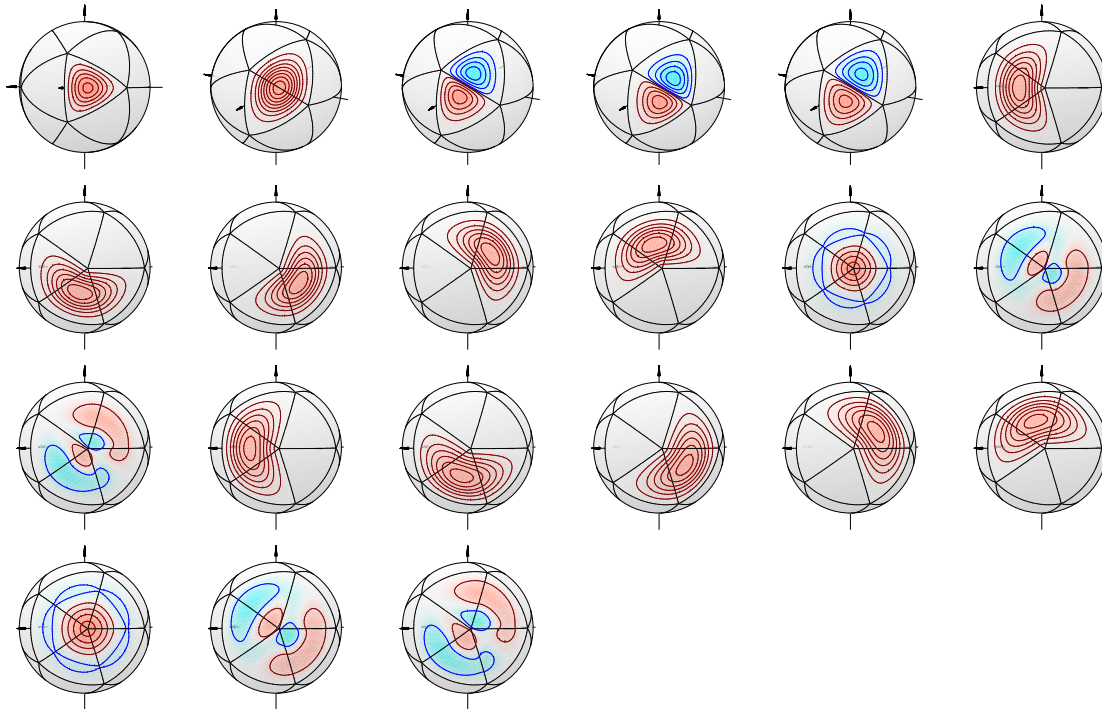
Original ANS basis

Original ANS basis for \mathcal{H}_1^7



Modified ANS basis

Modified ANS basis for $\mathcal{H}_1^6 \oplus \mathcal{H}_1^5$



Least Squares Approximation

Least squares approximation problem

Solve $Ac = b$ where

$$A_{i,j} = \langle \phi_i | \phi_j \rangle$$

$$b_i = \langle \phi_i | f \rangle$$

ϕ_1, \dots, ϕ_n = a function basis

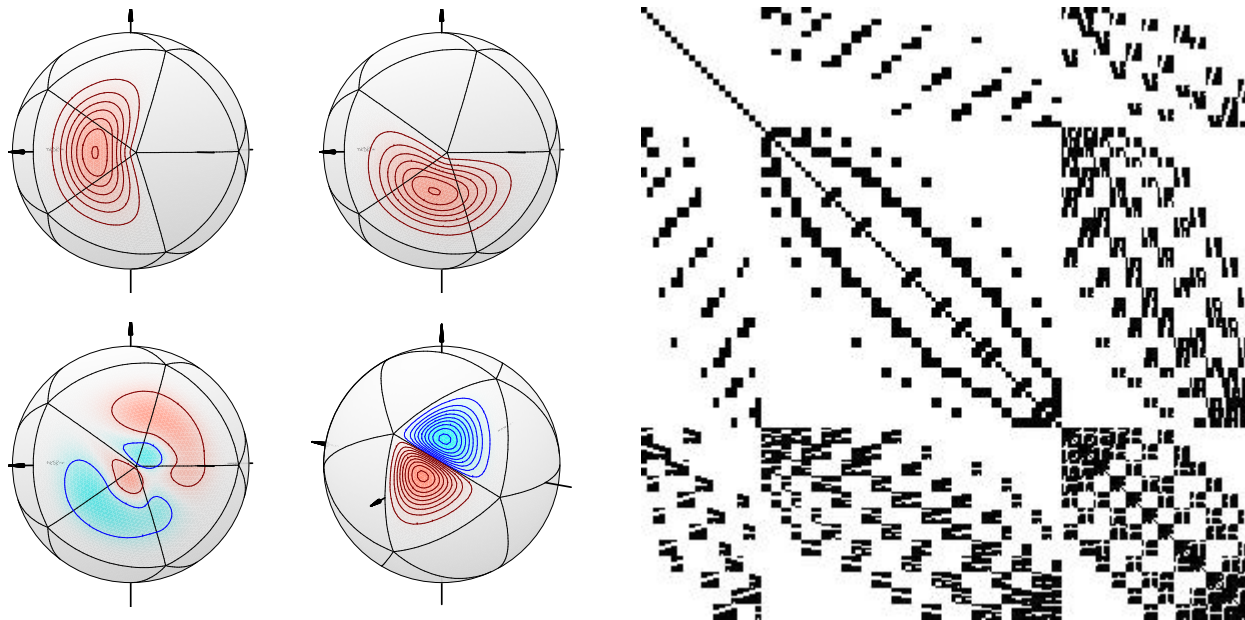
f = function to be approximated

$\langle \quad | \quad \rangle$ = a functional scalar product

c_1, \dots, c_n = coefficients of approximation

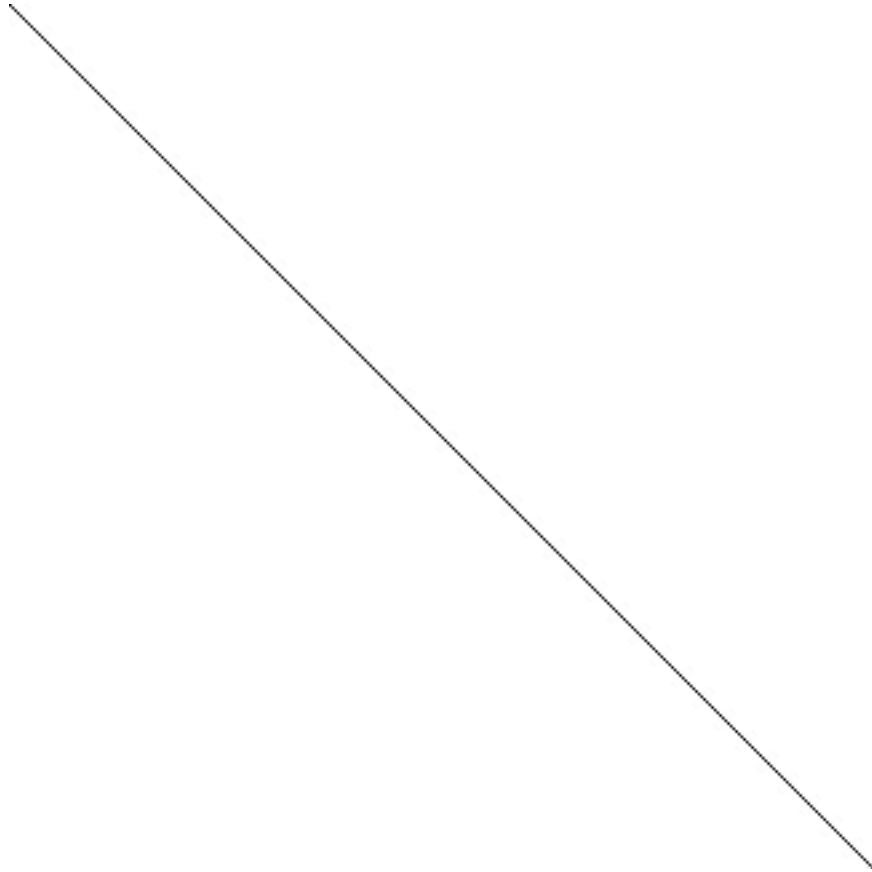
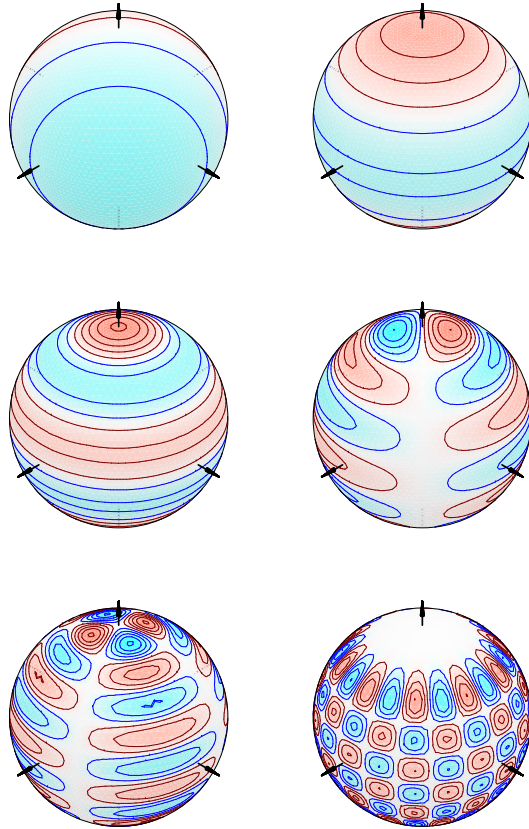
Matrices from Finite Elements

Finite-element bases give sparse matrices:



Orthogonal Bases

Orthogonal bases are much better:

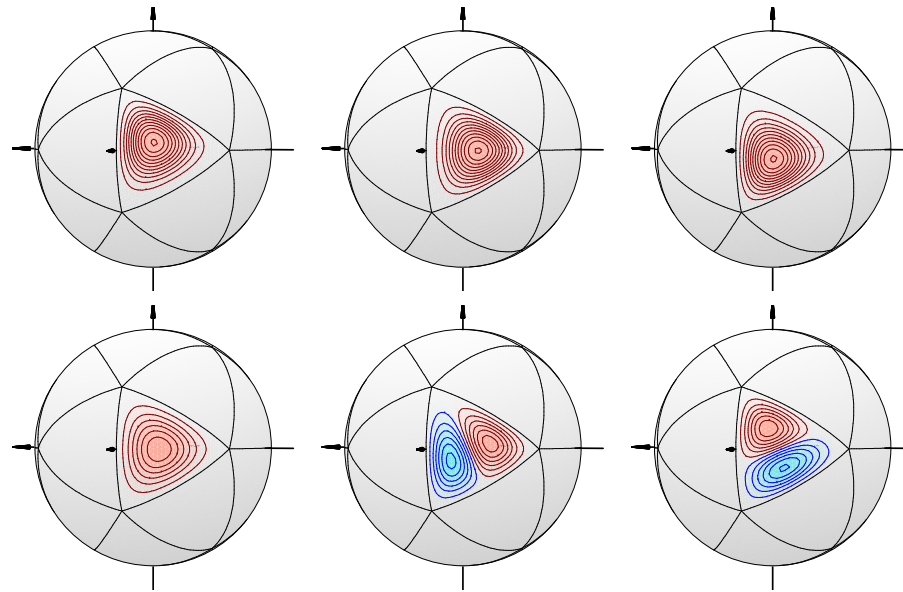


Semi-Orthogonalization

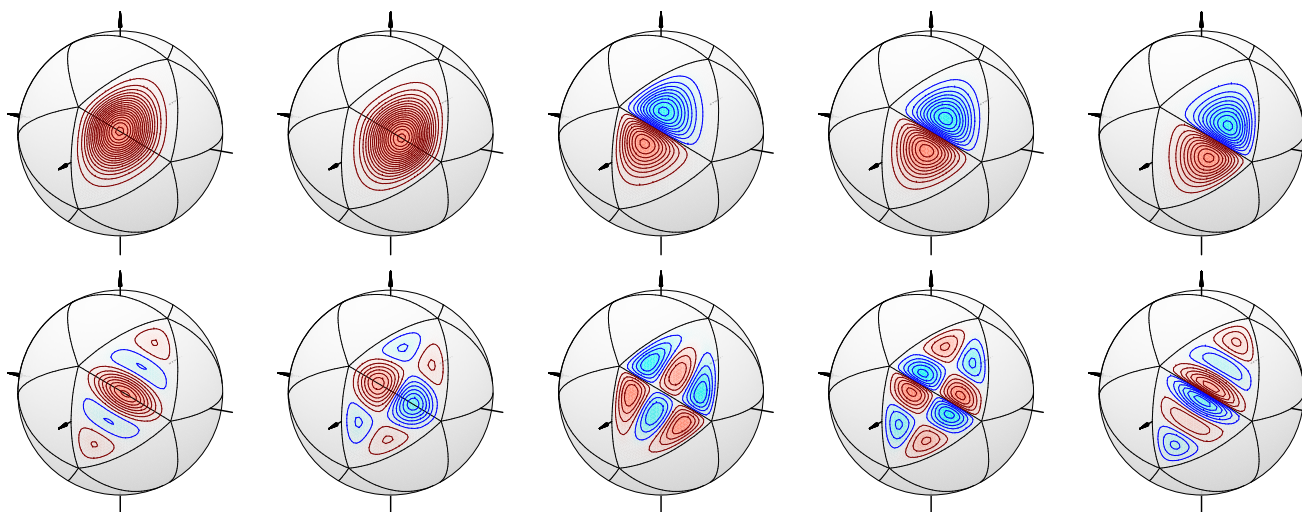
... but orthogonalization of a spline basis (e.g. Gram-Schmidt) destroys locality

Semi-orthogonalization:

Make the basis as orthogonal as possible
without increasing the element support



Semi-Orthogonalization (2)



Definitions and Notation

Ω = the domain.

T = the *mesh*: a partition of Ω into *tiles* t_1, \dots, t_n .

$\text{supp}(f)$ = *support* of f : set of tiles of T where f is not zero.

Φ = a *basis*: list of L.I. functions ϕ_1, \dots, ϕ_m on Ω

Φ_X = elements of Φ which have support X .

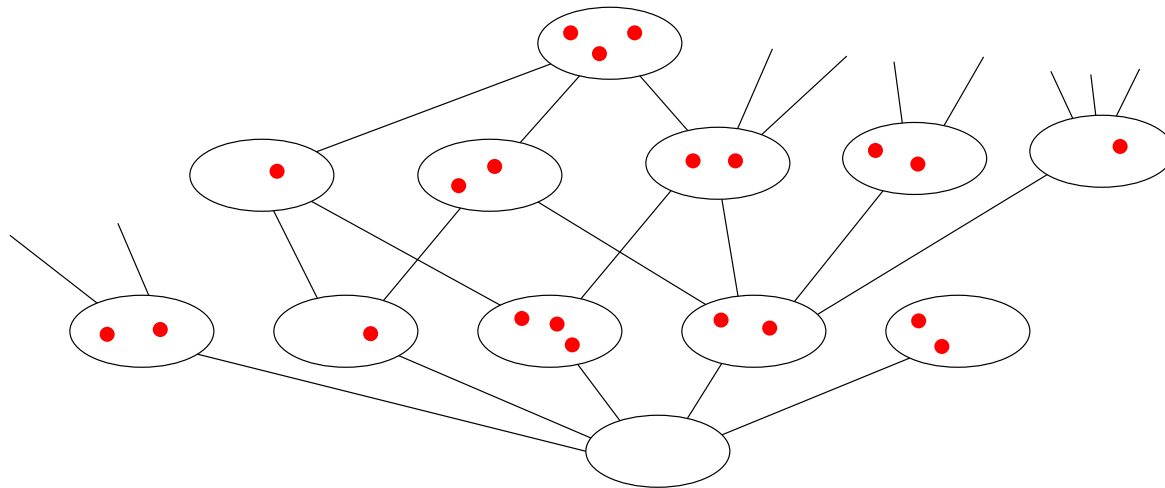
A basis Φ is *semi-orthogonal* iff
 $\langle \phi_i | \phi_j \rangle = 0$ whenever $\text{supp}(\phi_i) \subseteq \text{supp}(\phi_j)$.

Support Structure of a Spline Basis

The map $X \rightarrow \Phi_X$ defines the *support structure* of Φ

Viewed as a graph $\mathcal{G}(\Phi)$

- Nodes $\mathcal{V}\mathcal{G} = \{ X : X \subseteq T \}$.
- Edges $\mathcal{E}\mathcal{G}: X \rightarrow Y$ iff $X \subseteq Y$ and $|X| + 1 = |Y|$.
- Vertex X is decorated with the set Φ_X .



The useful part of $\mathcal{G}(\Phi)$ is small – $O(n) = O(m)$, not 2^m .

Semi-Orthogonalization Algorithm

- for each X in \mathcal{VG} , in increasing containment order
 - let $\Gamma_X \leftarrow \bigcup \{ \Phi_Y : Y \subset X \}$
 - for each element ϕ of Φ_X ,
 - make ϕ orthogonal to Γ_X
 - make the elements of Φ_X orthogonal among themselves.

Cost of Algorithm

Typically $|X|$, $|\Phi_X|$, and $|\Gamma_X|$ are $O(1)$ for any X with non-empty Φ_X .

No step is executed more than m times.

Total time is $O(m)$ (but with a respectable constant!).

Since \mathcal{VG} is processed bottom-up, Γ_X is semi-orthogonal.

If $\gamma_i \in \Gamma_X$ has $|\text{supp}(\gamma_i)| = 1$, we can assume that it is orthogonal to every other element in Γ_X .

When making $\Phi_X = (\gamma_1, \dots, \phi_k)$ orthogonal, use the eigenvectors of the matrix $A_X = \langle \gamma_i | \gamma_j \rangle$.

Example

Data used in test:

$\Omega =$ sphere

$M =$ icosahedral triangulation of \mathbf{S}^2

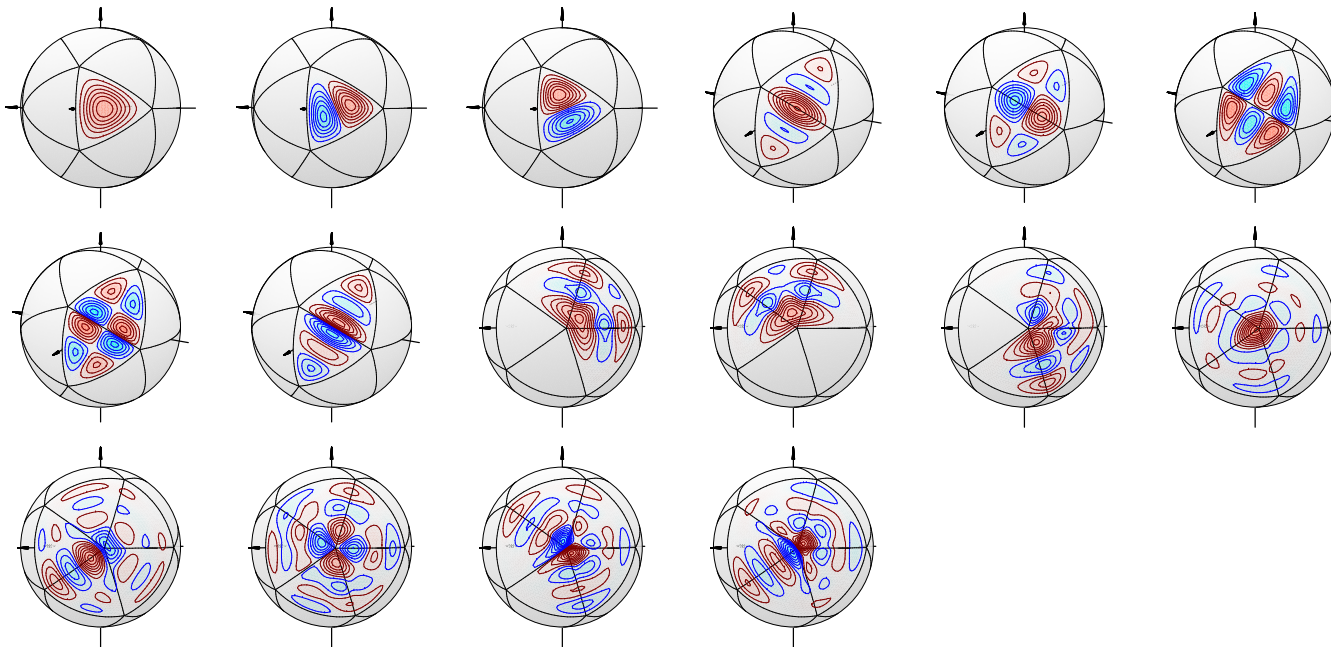
$\Phi =$ Alefeld-Neamtu-Schumaker (ANS) basis, \mathbf{C}_1

Two spaces:

- \mathcal{H}_1^7 - Homogeneous of degree 7 (dimension = 306)
- $\mathcal{H}_1^6 \oplus \mathcal{H}_1^5$ - Non-homogeneous of degree 6 (dimension = 332)

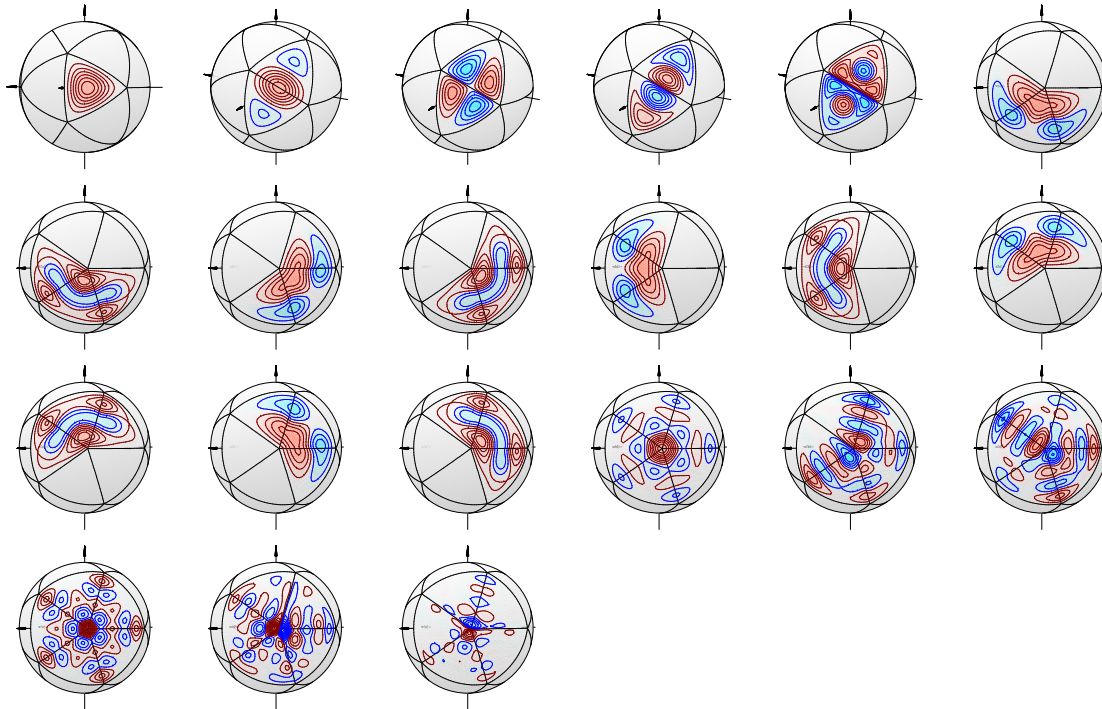
Original basis, semi-orthogonal

Semi-orthogonal ANS basis for \mathcal{H}_1^7



Modified basis, semi-orthogonal

Semi-orthogonal modified ANS basis for $\mathcal{H}_1^6 \oplus \mathcal{H}_1^5$



Comparison of bases

Non-zero elements and condition numbers

			plain		semi-ortho	
space	dim	basis	elems	cond	elems	cond
\mathcal{H}_1^7	306	O	21148	10^5	11316	10^2
$\overline{\mathcal{H}}_1^7$	306	N	18226	10^5	10662	10^2
$\mathcal{H}_1^6 + \mathcal{H}_1^5$	332	O	35232	10^{13}	29708	10^8
$\overline{\mathcal{H}}_1^6 + \overline{\mathcal{H}}_1^5$	332	N	29440	10^{13}	26515	10^8

Static Linear PDE on the Sphere

Consider the differential equation

$$(\mathcal{D}f)(p) = R(p)$$

for all $p \in \Omega$ where

- f is a function defined on Ω , to be determined.
- \mathcal{D} is a linear differential operator.
- R is a given function.

Example - Helmholtz equation

Helmholtz equation, $\Omega = \text{sphere}$

$$(\Delta f - cf)(p) = R(p)$$

where c is a constant,

Δ is the spherical Laplace-Beltrami operator, and

R is a given scalar field.

Aproximate Solution

Finite-element basis $\Phi = (\phi_0, \dots, \phi_{n-1})$.

Goal: compute

$$g = \sum_i c_i \phi_i$$

that approximates the solution f .

Galerkin's criterion:

$$\langle \mathcal{D}g - R \mid \phi_i \rangle = 0 \quad \text{for } i = 0, \dots, n - 1$$

where

$$\langle u \mid v \rangle = \int_{\Omega} u(p)v(p) dp.$$

Matrix Form

Galerkin's equations reduce to $Hc = b$ where

$$H_{ij} = \langle \mathcal{D}\phi_j - R \mid \phi_i \rangle$$

$$b_i = \langle R \mid \phi_i \rangle$$

For the Helmholtz equation

$$H_{ij} = \langle \Delta \phi_j \mid \phi_i \rangle - c \langle \phi_j \mid \phi_i \rangle$$

By Green's theorem, $\langle \Delta \phi_j \mid \phi_i \rangle = \langle \nabla \phi_j \mid \nabla \phi_i \rangle$
where ∇ is the spherical gradient

$$(\nabla f)(p) = (\text{grad } f)(p) - (p \langle \text{grad} \mid f \rangle)p$$

Non-Linear Problems

More generally

$$(\mathcal{D}f)(p) = R(f(p), p)$$

for all $p \in \Omega$, where R may depend non-linearly on f .

- solve by iteration.
- remove the linear part of R (Newton).

Some Results

RHS	sol	rms error (#iters)	
		\mathcal{H}^7	$\mathcal{H}^6 \oplus \mathcal{H}^5$
sqr $6.5x^2 - 2$	x^2	4.5×10^{-5}	1.0×10^{-12}
cub $12.5x^3 - 6x$	x^3	9.0×10^{-13}	1.0×10^{-12}
qrt $20.5x^4 - 12x^2$	x^4	4.4×10^{-5}	2.0×10^{-12}
sep $56.5x^7 - 42x^5$	x^7	9.0×10^{-13}	6.9×10^{-5}
oct $72.5x^8 - 56x^6$	x^8	8.9×10^{-5}	8.6×10^{-5}
cos $(1.5 - x^2) \cos x - 2x \sin x$	$\cos x$	8.1×10^{-5}	2.1×10^{-9}
sin $(1.5 - x^2) \sin x + 2x \cos x$	$\sin x$	3.1×10^{-10}	1.3×10^{-8}
exp $e^x(x^2 + 2x - 0.5)$	e^x	1.1×10^{-4}	1.4×10^{-8}
mcos $(1.5 - x^2)(3 \cos x + f)/4 - 2x \sin x$	$\cos x$	8.1×10^{-5} (25)	2.6×10^{-9} (40)
mexp $0.5(e^x + f)(x^2 + 2x - 0.5)$	e^x	1.1×10^{-4} (19)	1.4×10^{-8} (35)

Dynamic Problems

Time-dependent problem

$$(\mathcal{D}f)(p, t) = R(f(p, t), p, t)$$

where

p = point of Ω

t = time

f = function to be determined

\mathcal{D} = a linear differential operator

R = a given function.

Heat Diffusion on Rotating Sphere

Example - Heat diffusion on a rotating sphere

$$\left(\frac{\partial}{\partial t}f\right)(p, t) - K (\Delta f)(p, t) + \omega V(p) \cdot (\nabla f)(p, t) + L f(p, t) = R(p, t)$$

where

K = heat diffusion coefficient

ω = angular speed

V = velocity field for $\omega = 1$: $V(x, y, z) = (-y, x, 0)$

L = coefficient of linear heat loss

R = external heat input (may depend on $f(p, t)$)

Petrov-Galerkin Approach

Assume a basis of space-time elements

$$\Psi = (\psi_0, \dots, \psi_{N-1})$$

where each ψ_ℓ is a function of $\Omega \times \mathbf{R}$.

We look for an approximation h to f

$$h(p) = \sum_{\ell} c_{\ell} \psi_{\ell}(p, t)$$

where the c_{ℓ} are coefficients to be determined.

Petrov-Galerkin Approach (2)

We can use Petrov-Galerkin by treating time as another space dimension:

$$\langle \mathcal{D}h - R \mid \theta_k \rangle = 0$$

for all k , where

$$\Theta = (\theta_0, \dots, \theta_{N-1})$$

is another basis of space-time elements (*gauge functions*).

We get the non-linear system

$$Mc = b$$

where

$$M_{k\ell} = \langle \mathcal{D}\psi_\ell \mid \theta_k \rangle \quad b_k = \langle R \mid \theta_k \rangle$$

Separating Time and Space

We can partially separate the variables by using tensor-type bases for Φ and Θ . Let

$$\Phi = (\phi_0, \dots, \phi_{n-1}) = \text{a basis over } \Omega$$

$$\Lambda = (\lambda_0, \dots, \lambda_{m-1}) = \text{a basis over } \mathbf{R}$$

$$\Gamma = (\gamma_0, \dots, \gamma_{m-1}) = \text{another basis over } \mathbf{R}$$

Then use

$$\psi_k = \phi_i \lambda_r$$

$$\theta_k = \phi_i \gamma_r$$

where $i = 0, \dots, n - 1$, $r = 0, \dots, m - 1$, and $k = in + r$.

Separating the Integrals

Integrals can be separated:

$$\langle \psi_\ell | \theta_k \rangle = \langle \phi_j | \phi_i \rangle \langle \lambda_s | \gamma_r \rangle$$

$$\langle \Delta \psi_\ell | \theta_k \rangle = \langle \Delta \phi_j | \phi_i \rangle \langle \lambda_s | \gamma_r \rangle$$

$$\langle V \cdot \nabla \psi_\ell | \theta_k \rangle = \langle V \cdot \nabla \phi_j | \phi_i \rangle \langle \lambda_s | \gamma_r \rangle$$

$$\langle \frac{\partial}{\partial t} \psi_\ell | \theta_k \rangle = \langle \phi_j | \phi_i \rangle \langle \frac{\partial}{\partial t} \lambda_s | \gamma_r \rangle$$

Progressive Solution

Specifically, for the sphere we use

Φ = the ANS spherical splines C_μ

Λ = C_μ spline pulses spanning 2 time steps

Γ = spline pulses of degree $\leq \mu$ spanning one step

The time element λ_r is centered at epoch $\lfloor r/q \rfloor$ where $q = \mu + 1$.

Progressive Solution (2)

The system's matrix M then has the block structure

$$\begin{pmatrix} N_1 & N_0 & 0 & 0 & 0 & 0 & \cdots \\ \vdots & \ddots & \cdots & \cdots & \cdots & \cdots & \cdots \\ \vdots & 0 & N_1 & N_0 & 0 & 0 & \cdots \\ \vdots & 0 & 0 & N_1 & N_0 & 0 & \cdots \\ \vdots & 0 & \vdots & 0 & N_1 & N_0 & \cdots \\ \vdots & \ddots & \vdots & \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$

where each block N_k , for $k=0..1$, is the a matrix of the scalar products of the differentiated basis elements in epoch $j - k$ against the gauge functions in epoch j .

Then at each step we solve

$$N_0 a_0 == d - N_1 a_1$$

where a_0 is the coefficients for the present epoch, and a_1 those of the previous epoch, and d is the corresponding segment of b .

An Example Result

